Investigation of the Impact of Ocean Waves on Development of Explosive Cyclone with a Coupled Model

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Granted by JSPS Kakenhi-17J07627, Kakenhi-16H01846

Interactions among Atmosphere, Ocean, & Wave

Ocean - Atmosphere

- SST feedback to heat flux
- Current feedback to momentum
- Water flux takes rainfall into account

Atmosphere - Wave

- Surface roughness
- Sea spray affects water and heat

Ocean - Wave

- Wave-induced current and turbulence (Stokes drift, Langmuir circulation)
- Enhanced mixing due to wave breaking
- Wave setup
- Wave-induced pressure



Explosive Cyclone (EC)

- An explosive cyclone is generated in the area of high baroclinicity in winter and spring. Its definition is that its central SLP drops over 24hPa within a day. The explosive cyclone is hard to predict accurately and causes sever disasters in broad areas.
- Marine disasters caused by ECs
 - Coastal facility damages (Toyama, 2008)
 - Downfall and stranding of vessels (Ibaraki, 2006)
- Main factors of the development
 - Large anomaly of potential vorticity in upper layer
 - Inflow of latent heat and vapor from ocean
 - Latent heat release from cloud and rain generation







Influence of Kuroshio on explosive cyclone (Nonaka et al., 2016)

Research Overview

Motivation

• The Atmospheric and ocean model cannot resolve ocean wave effect explicitly.

• The effects of ocean wave on an meso-scale phenomenon in atmosphere and ocean are remain to be clarified for the sake of accurate weather forecast.

Objective

To clarify ocean wave effects on the boundary layer turbulence in both atmosphere and ocean and on the development of a winter explosive cyclone.

Methodology

 An atmosphere-ocean-wave coupled model simulates an explosive cyclone considering ocean wave effects.

Explosive Cyclone in Northwestern Atlantic in 2018 Jan

In 2018 January, an explosive cyclone emerged in Northwestern Atlantic and developed as much as 949hPa in its central SLP. It caused enormous blizzard and damages of \$11 billion.





NOAA GOES-16 satellite

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Configuration of the Coupled Model

	WRF	WAVEWATCH III	CROCO
Domain	la	on: $30^{\circ} \sim 50^{\circ}$, $lat: -88^{\circ} \sim -50^{\circ}$	
Period		2018/01/01 ~ 2018/01/10	
Resolution	°80.0	0.08°	1/18°
Time Step	30s	360s (global)	9s(2D), 360s(3D)
Grid	lon 501 × lat 326,60 layer	lon 536 × lat 349	lon 729 × lat 597,45layer
Others	Initial/boundary: CFSR Boundary layer: MYNN scheme	Initial/boundary: Atlantic Model wind, ice forcing: CFSR	Initial/boundary: HYCOM No tide, σ coordinate



Configuration of Surface Physics

Surface Stress: $ au_{tot} = \rho_a u_*^2 = \rho_a C_d U^2$
Surface Sensible Heat Flux: $H_S = -\rho_a c_p u_* \theta_* = -\rho_a c_p C_h U(\theta_a - \theta_s)$
Surface Latent Heat Flux: $H_L = L_e \rho_a u_* q_* = L_e \rho_a M C_q U (q_s - q_a)$
Charnock Parameter: $\alpha = 0.01(1 - \tau_w/\tau_{tot})^{-1/2}$ (Janssen 1991)
Surface Roughness: $z_0 = z_{0\nu} + z_w = \frac{0.11\nu}{u_*} + \alpha \frac{u_*^2}{g} = \frac{0.11\nu}{u_*} + \frac{0.01}{\sqrt{1 - \tau_w / \tau_{tot}}} \frac{u_*^2}{g}$ ($z_{0,max} = 2.85mm$; Davis et al. 2008)
Wave-induced Stress: $\tau_w = \rho_w \int_0^\infty \int_{-\pi}^{\pi} f^2 S_{in} \frac{\vec{k}}{k} d\theta df$
Friction Velocity: $u_{*a} = kU[\ln(z/z_0) - \psi_m(z/L)], \ u_{*w} = u_{*a}(1 + \gamma w_{age}[1 - \cos\phi])$ ϕ : angle between wind – wave, $\gamma = 0.007, w_{age} = C_p/u_{*a}$
(Patton et al., 2019)

Surface roughness reflects wave development Friction velocity reflects misalignment of wave-wind direction

Governing Equations of CROCO (new ROMS-AGRIF)

- Make use of the wave-current Interactions configuration in Uchiyama et al. (2010).
- Stokes Drift: (u^{st}, v^{st}, w^{st}) , Vortex force: (V_x, V_y, V_z) , Bernoulli head: ϕ_B , Unpreserved force: *F* Reynolds-averaged governing equations are:

$$\text{Momentum:} \begin{array}{l} \frac{\partial u}{\partial t} + u_j \frac{\partial u}{\partial x_j} - fv = -\frac{\partial(\phi + \phi_B)}{\partial x} - \frac{\partial}{\partial x_j} \left(\overline{u'_j u} - v \frac{\partial u}{\partial x_j} \right) + V_x + F_x + F_x^w \\ \frac{\partial v}{\partial t} + u_j \frac{\partial v}{\partial x_j} + fu = -\frac{\partial(\phi + \phi_B)}{\partial y} - \frac{\partial}{\partial x_j} \left(\overline{u'_j v} - v \frac{\partial v}{\partial x_j} \right) + V_y + F_y + F_y^w \\ \text{Scalor:} \quad \frac{\partial c}{\partial t} + u_j \frac{\partial c}{\partial x_j} = -u_j^{st} \frac{\partial c}{\partial x_j} - \frac{\partial}{\partial x_j} \left(\overline{u'_j c} - v_c \frac{\partial c}{\partial x_j} \right) + F_c \quad (T, S) \\ \text{Pressure:} \quad \frac{\partial \phi}{\partial z} + \frac{g\rho}{\rho_0} = -\frac{\partial \phi_B}{\partial z} + V_z \\ \text{Continuity:} \quad \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \\ \left(V_x, V_y \right) = -\hat{\mathbf{z}} \times \mathbf{u}^{\text{st}} [(\hat{\mathbf{z}} \cdot \nabla_\perp \times \mathbf{u}) + f] - w^{st} \frac{\partial \mathbf{u}}{\partial z}, \\ \phi_B = \frac{\mathbf{u}_z}{4 \frac{1}{k \sinh^2[kD]}} \int_{-h}^z \frac{\partial(\mathbf{k} \cdot \mathbf{u})^2}{\partial \zeta^2} \sinh[2k(z-\zeta)] d\zeta \\ \text{Zoin the transformational Workshop on Waves, Storm Surges and Coastal Hazards} \end{array}$$

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Turbulence Parameterization: k-kl model (MY25)

- In the CROCO model, several GSL (generic length scale) approaches are provided (Umlauf and Burchard, 2003): e.g., k kl, k ε, k ω model. They are two-equation models which diagnose two turbulence values explicitly.
- Here, Mellor–Yamada level 2.5 (MY25, k kl) scheme is selected, which is originally proposed in Mellor and Yamada (1974).

For turbulent kinetic energy: $\mathbf{E} = \frac{1}{2} \left(\overline{u'^2} + \overline{v'^2} + \overline{w'^2} \right)$ and turbulent length scale: \mathbf{l} ,

$$\frac{\partial E}{\partial t} + u_j \frac{\partial E}{\partial x_j} = \frac{\partial}{\partial z} \left(K_Q \frac{\partial E}{\partial z} \right) + P + B - \varepsilon$$

$$\frac{\partial(El)}{\partial t} + u_j \frac{\partial(El)}{\partial x_j} = \frac{\partial}{\partial z} \left[K_Q \frac{\partial(El)}{\partial z} \right] + l(c_1 P + c_3 B - c_2 \varepsilon F_{wall})$$

shear production: $P = -\overline{u'w'}\frac{\partial u}{\partial z} - \overline{v'w'}\frac{\partial v}{\partial z}$, buoyancy production: $B = -\frac{g}{\rho_0}\overline{\rho'w'}$, dissipation: ε $K_Q = \sqrt{2k}lS_q$; $S_q = 0.2$; $c_1 = 0.9, c_3 = 0.9, c_2 = 0.5$

Parameter used here is proposed by Umlauf et al. (2003) and Canuto et al. (2001).

Wave-induced Mixing Parameter in MY25 scheme

Babanin (2011) suggested non-breaking wave-induced mixing formula, like: shear production: $P \rightarrow P + P_w = -\overline{u'w'}\frac{\partial u}{\partial z} - \overline{v'w'}\frac{\partial v}{\partial z} + bk\left(\frac{\omega H_s}{2}e^{kz}\right)^3$

b = 0.0014, ω : angular frequency, H_s : significant wave height, k: wave number

Aijaz et al. (2017) examined this P_w impacts under a tropical cyclone condition.

Experiment Conditions

1. Examine Atmosphere-Wave Interactions

2. Examine Ocean-Wave Interactions

without wave-atmosphere coupling

3. Examine Atmosphere-Ocean-Wave Interactions

Overview of Wave Coupling Effects

Overview of the Simulation Results

Surface Wind (10m height)

Surface Wind Speed (A,MYNN)

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Wind Speed at 10m height (m/s) at True U at 10 M (m s-1) at True Sea Level Pressure (hPa) at True

Significant Wave Height

Atmosphere-Wave Coupling Effects (AW – A)

Friction Velocity u_*

 u_* and z_0 are intensified by the modified parameterization in the atmosphere-wave coupling

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Surface Roughness z_0

Atmosphere-Wave Coupling Effects (AW – A)

Surface Heat Flux

Sea Level Pressure (SLP)

Surface heat flux raises with modified z_0 and u_* , and SLP drops more deeply in AW

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55°W

60°W

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Surface Parameters: AW, A

- Atmosphere-wave coupling makes surface rougher and raises friction velocity.
- Ocean wave coupling increases drag coefficient and enthalpy coefficient, resulting in higher heat flux into the atmosphere.

✓ These results are consistent with Pianezze et al. (2018)

Ocean-Wave Coupling Effects (AWO. naw – AO)

* Defined as the depth whose buoyancy frequency is the largest

Sea Surface Temperature (SST)

Mixed Layer Depth (MLD)*

0

16

32

-32

-16

With wave-induced mixing, MLD and SST drops more than independent ocean model.

With wave-induced mixing, the eddy viscosity is amplified and mixed layer depth is deepened.

Area-Averaged Sea Level Pressure & Heat Flux (AO, AWO. naw)

Atmosphere-Ocean-Wave Coupling Effects (AWO – AO)

With atmosphere-ocean-wave interactions, SST and MLD drops more deeply.

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50°N

45°N

40°N

35°N

30°N

25°N

60°W

16

55°W

32

Atmosphere-Ocean-Wave Coupling Effects (AW0 – A0)

Surface Heat Flux

Sea Level Pressure (SLP)

0

In fully coupled model, the surface hear flux is enhanced and SLP drops more deeply.

Area-Averaged Sea Level Pressure & Heat Flux (AO, AWO)

Ocean Wave Effects on the Development of Explosive Cyclone

Summary

- The modification of z_0 and u_* increases surface stress and roughness of the atmospheric surface, and enhances the surface heat flux.
- Introducing the wave-induced mixing in the ocean model declines MLD more deeply and cools SST down.
- The wave coupling to both atmosphere and ocean intensifies the inertial oscillations by the explosive cyclone, resulting in deeper MLD and cooler SST.
 These wave coupling effects influence the surface heat flux significantly and make prominent changes in the development of the explosive cyclone.

Future work

It remains to be clarified how the enhanced heat flux make changes in the development in the context of atmospheric dynamism.

Feedbacks of the wave coupling to the wave development are still unclear.

APPENDIX

Turbulence Parameterization: KPP (LMD scheme)

KPP (K-Profile Parameterization)

• Parameterize 2nd-order fluxes with 1st-order value like:

$$\overline{u'w'} = -K_M \frac{\partial u}{\partial z}, \qquad \overline{v'w'} = -K_M \frac{\partial v}{\partial z}, \qquad \overline{w'\theta'} = -K_H \frac{\partial \theta}{\partial z}$$

• In the KPP scheme, non-local term γ_a is added to:

$$\overline{u'w'} = -K_M \left(\frac{\partial u}{\partial z} - \gamma_u\right), \quad \overline{v'w'} = -K_M \left(\frac{\partial v}{\partial z} - \gamma_v\right), \quad \overline{w'\theta'} = -K_H \left(\frac{\partial \theta}{\partial z} - \gamma_\theta\right)$$

In the CROCO model, LMD (Large, McWilliams, Doney 1994) scheme is used.

In terms of wave-current interactions, Qiao et al. (2004) introduced nonbreaking wave-induced turbulence effect into the KPP scheme.

New
$$K_M \to K_M + B_v$$
, New $K_H \to K_H + B_v$
, where $B_v = \alpha A^3 k \omega \left[\frac{\sinh k(h+z)}{\sinh kh} \right]^3$ ($\alpha = 0.1$; cf. Wang et al. 2010)

Comparison with Marine Observations

Wind Speed at LaHave Bank (2018/01) Wind[m/s] WRF AW Buoy SLP at LaHave Bank (2018/01) WRF AW Buoy SLP[hPa] Hs at LaHave Bank (2018/01) WRF AW Buoy Wave height[m] Date

Pressure tendency equation (Fink et al., 2012)

$$\frac{\partial p_{sfc}}{\partial t} = \rho_{sfc} \frac{\partial \phi_{p_2}}{\partial t} + \rho_{sfc} R_d \int_{p_{sfc}}^{p_2} \frac{\partial T_v}{\partial t} d\ln p + g(E - P) + RES_{PTE}$$
Geopotential tendency ITT: integrated virtual temperature tendency
$$\rho_{sfc} R_d \int_{p_{sfc}}^{p_2} \frac{\partial T_v}{\partial t} d\ln p = \rho_{sfc} R_d \int_{p_{sfc}}^{p_2} \left(-\vec{v} \cdot \vec{\nabla_p} T_v \right) d\ln p$$
TADV: horizontal temperature advection
$$+ \rho_{sfc} R_d \int_{p_{sfc}}^{p_2} \left(\frac{R_d T_v}{c_p p} - \frac{\partial T_v}{\partial p} \right) \omega d\ln p$$
VMT: vertical motions
$$+ \rho_{sfc} R_d \int_{p_{sfc}}^{p_2} \frac{T_v Q}{c_p T} d\ln p$$
DIAB: diabatic processes
$$+ RES_{ITT}$$

ITT: integrated virtual temperature tendency

AW-A: MSLP Difference (MYNN)

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AW-A: ITT (MYNN)

Init: 2018-01-01_00:00:00 Valid: 2018-01-04_18:00:00

vertically integrated virtual temperature tendency diff (hPa/6hour) at True Sea Level Pressure (hPa) at True

