



# Coherent Interference and Diffraction in Random Waves

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# Outline

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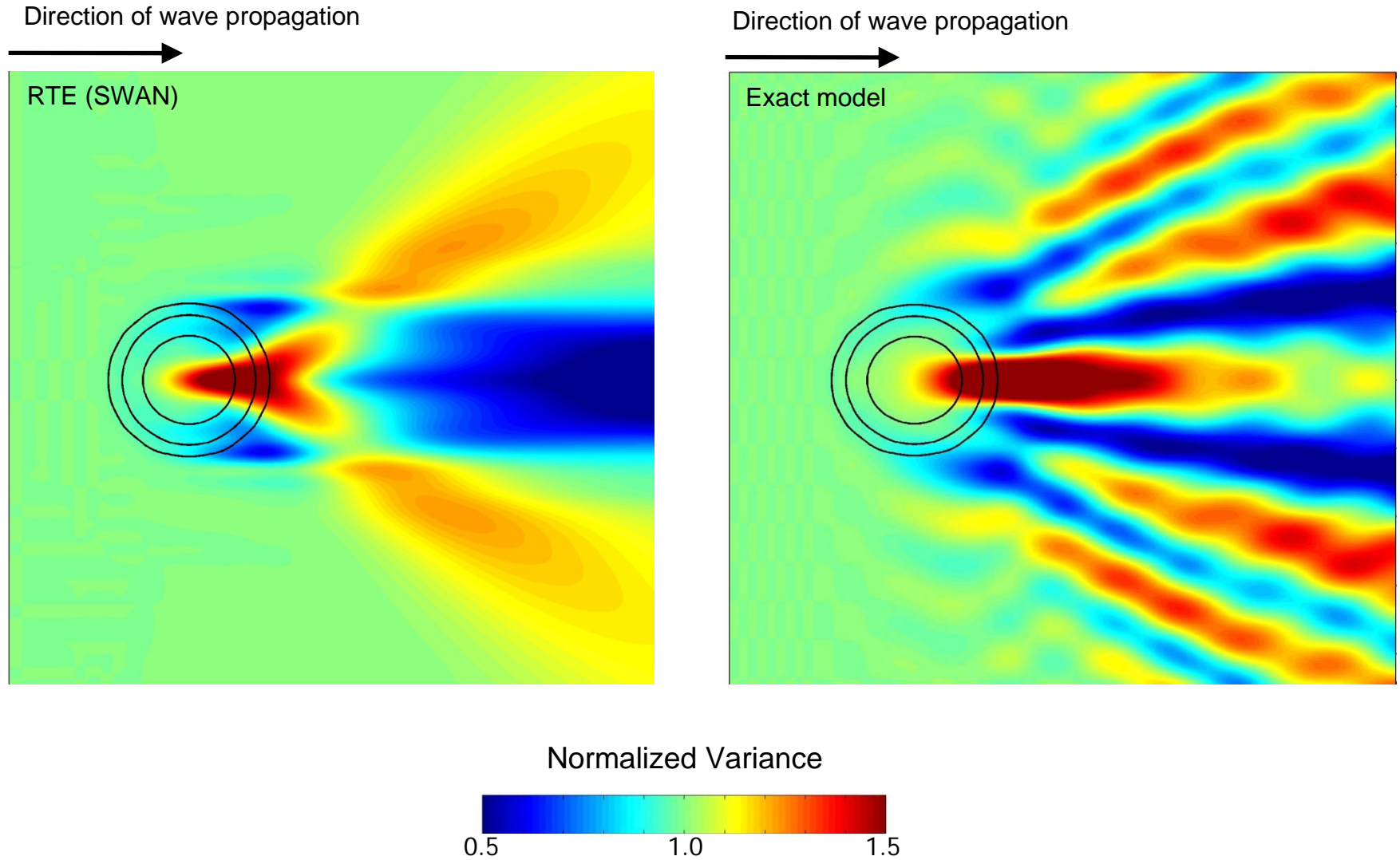
- What is wrong with the RTE?
- How do we improve it?
- Examples
- Concluding remarks

# Radiative transport equation (RTE)

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$$\frac{\partial \mathcal{E}}{\partial t} + \nabla_{\mathbf{k}} \Omega \cdot \nabla_{\mathbf{x}} \mathcal{E} - \nabla_{\mathbf{x}} \Omega \cdot \nabla_{\mathbf{k}} \mathcal{E} = S$$

# What is wrong with the RTE`?



# What is wrong with the RTE?

Consider a two-component wavefield described by:

$$\zeta(\mathbf{x}, t) = \hat{\zeta}_1 \exp [i\Phi_1] + \hat{\zeta}_2 \exp [i\Phi_2] \quad \Phi_j = \mathbf{k}_j \cdot \mathbf{x} - \omega_j t$$

Variance:

$$\langle \zeta \zeta^* \rangle = \underbrace{\langle \hat{\zeta}_1 \hat{\zeta}_1^* \rangle}_{\text{auto-variance}} + \underbrace{\langle \hat{\zeta}_2 \hat{\zeta}_2^* \rangle}_{\text{auto-variance}} + \underbrace{\left( \langle \hat{\zeta}_1 \hat{\zeta}_2^* \rangle \exp [i\Phi_1 - i\Phi_2] + C.C. \right)}_{\text{cross-variance}}$$

auto-variance  
contributions

cross-variance  
contributions

How do transport these interference terms?

# Evolution of interference contributions

Consider the transform pair for the wave variable  $\zeta$  :

$$\zeta(\mathbf{x}, t) = \int \widehat{\zeta}(\mathbf{k}, t) \exp(i\mathbf{k} \cdot \mathbf{x}) d\mathbf{k},$$

$$\widehat{\zeta}(\mathbf{k}, t) = \frac{1}{(2\pi)^2} \int \zeta(\mathbf{x}, t) \exp(-i\mathbf{k} \cdot \mathbf{x}) d\mathbf{x}.$$

Local dispersion relation in slowly changing medium:

$$\omega = \Omega(\mathbf{k}, \mathbf{x}, t).$$

Using operator correspondence

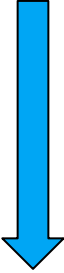
$$-i\omega \rightarrow \partial_t, \quad \mathbf{x} \rightarrow i\nabla_{\mathbf{k}},$$

$$\partial_t \widehat{\zeta}(\mathbf{k}, t) = -i\Omega(\mathbf{k}, i\nabla_{\mathbf{k}}, t) \widehat{\zeta}(\mathbf{k}, t)$$

# Evolution of interference contributions

A transport equation for  $\gamma(\mathbf{k}_1, \mathbf{k}_2, t) = \langle \widehat{\zeta}(\mathbf{k}_1, t) \widehat{\zeta}^*(\mathbf{k}_2, t) \rangle$

$$[\partial_t + i\Omega(\mathbf{k}_1, i\nabla_{\mathbf{k}_1}, t) - i\Omega(\mathbf{k}_2, -i\nabla_{\mathbf{k}_2}, t)] \gamma(\mathbf{k}_1, \mathbf{k}_2, t) = 0$$

  $\left\{ \begin{array}{l} - \text{Introduce: } \mathbf{k} = \frac{1}{2}(\mathbf{k}_1 + \mathbf{k}_2), \quad \mathbf{u} = \mathbf{k}_1 - \mathbf{k}_2, \\ - \text{take the inverse } \mathbf{u} \rightarrow \mathbf{x} \quad \text{transform} \end{array} \right.$

$$\partial_t \mathcal{E} = -i \left[ \Omega \left( \mathbf{k} - \frac{i}{2} \nabla_{\mathbf{x}, \mathbf{x}} + \frac{i}{2} \nabla_{\mathbf{k}, t} \right) - \Omega \left( \mathbf{k} + \frac{i}{2} \nabla_{\mathbf{x}, \mathbf{x}} - \frac{i}{2} \nabla_{\mathbf{k}, t} \right) \right] \mathcal{E}$$

where

$$\mathcal{E}(\mathbf{k}, \mathbf{x}, t) = \frac{1}{2} \int \left\langle \widehat{\zeta} \left( \mathbf{k} + \frac{\mathbf{u}}{2}, t \right) \widehat{\zeta}^* \left( \mathbf{k} - \frac{\mathbf{u}}{2}, t \right) \right\rangle \exp(i\mathbf{u} \cdot \mathbf{x}) d\mathbf{u}$$



# The Coupled Mode Spectrum

Consider  $\mathcal{E}(\mathbf{k}, \mathbf{x}, t)$  (Coupled Mode Spectrum)

- Transform of the two point correlator:  $\left\langle \hat{\zeta}\left(\mathbf{k} + \frac{\mathbf{u}}{2}, t\right) \hat{\zeta}^*\left(\mathbf{k} - \frac{\mathbf{u}}{2}, t\right) \right\rangle$
- Marginal relation  $\mathcal{V}(\mathbf{x}, t) = \int \mathcal{E}(\mathbf{k}, \mathbf{x}, t) d\mathbf{k}$ . represents Variance

However, it is not a variance density spectrum

- Captures complete second order statistics
- Is real, but not point-wise positive

# Evolution of interference contributions

Governing equation

$$\partial_t \mathcal{E} = -i \left[ \underbrace{\Omega \left( \mathbf{k} - \frac{i}{2} \nabla_{\mathbf{x}}, \mathbf{x} + \frac{i}{2} \nabla_{\mathbf{k}}, t \right)}_{\text{Taylor series for the operators}} - \underbrace{\Omega \left( \mathbf{k} + \frac{i}{2} \nabla_{\mathbf{x}}, \mathbf{x} - \frac{i}{2} \nabla_{\mathbf{k}}, t \right)}_{\text{Taylor series for the operators}} \right] \mathcal{E}$$



Taylor series for the operators

$$\frac{\partial \mathcal{E}}{\partial t} + \sum_{n=1}^{\infty} \left[ \sum_{i_1=1}^4 \cdots \sum_{i_{2n-1}=1}^4 \mathcal{C}_{i_1, \dots, i_{2n-1}} \left( \prod_{m=1}^{2n-1} \frac{\partial}{\partial q_{i_m}} \right) \right] \mathcal{E} = 0$$



First order...

$$\boxed{\frac{\partial \mathcal{E}}{\partial t} + \nabla_{\mathbf{k}} \Omega \cdot \nabla_{\mathbf{x}} \mathcal{E} - \nabla_{\mathbf{x}} \Omega \cdot \nabla_{\mathbf{k}} \mathcal{E} = 0}$$

Radiative transport equation!

# Quasi-Coherent Approximation

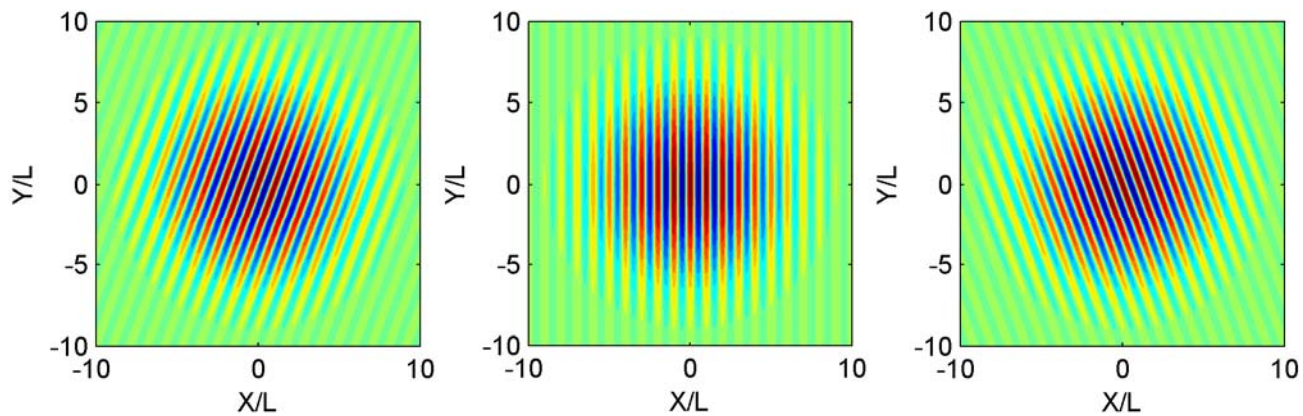
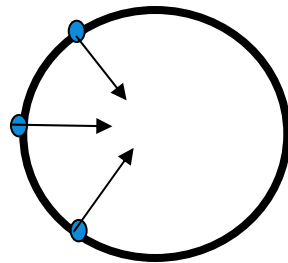
The Quasi-Coherent (QC) approximation (with  $\mathbf{q} = [\mathbf{x}, \mathbf{k}]^T$ )

$$\left[ \underbrace{\frac{\partial}{\partial t} + c_i \frac{\partial}{\partial q_i}}_{\text{Rad. Transp. Equation}} + \underbrace{c_{i,j,k} \frac{\partial^3}{\partial q_i \partial q_j \partial q_k}}_{\text{Corrections for coherent interference}} \right] \mathcal{E} = 0$$

Rad. Transp. Equation

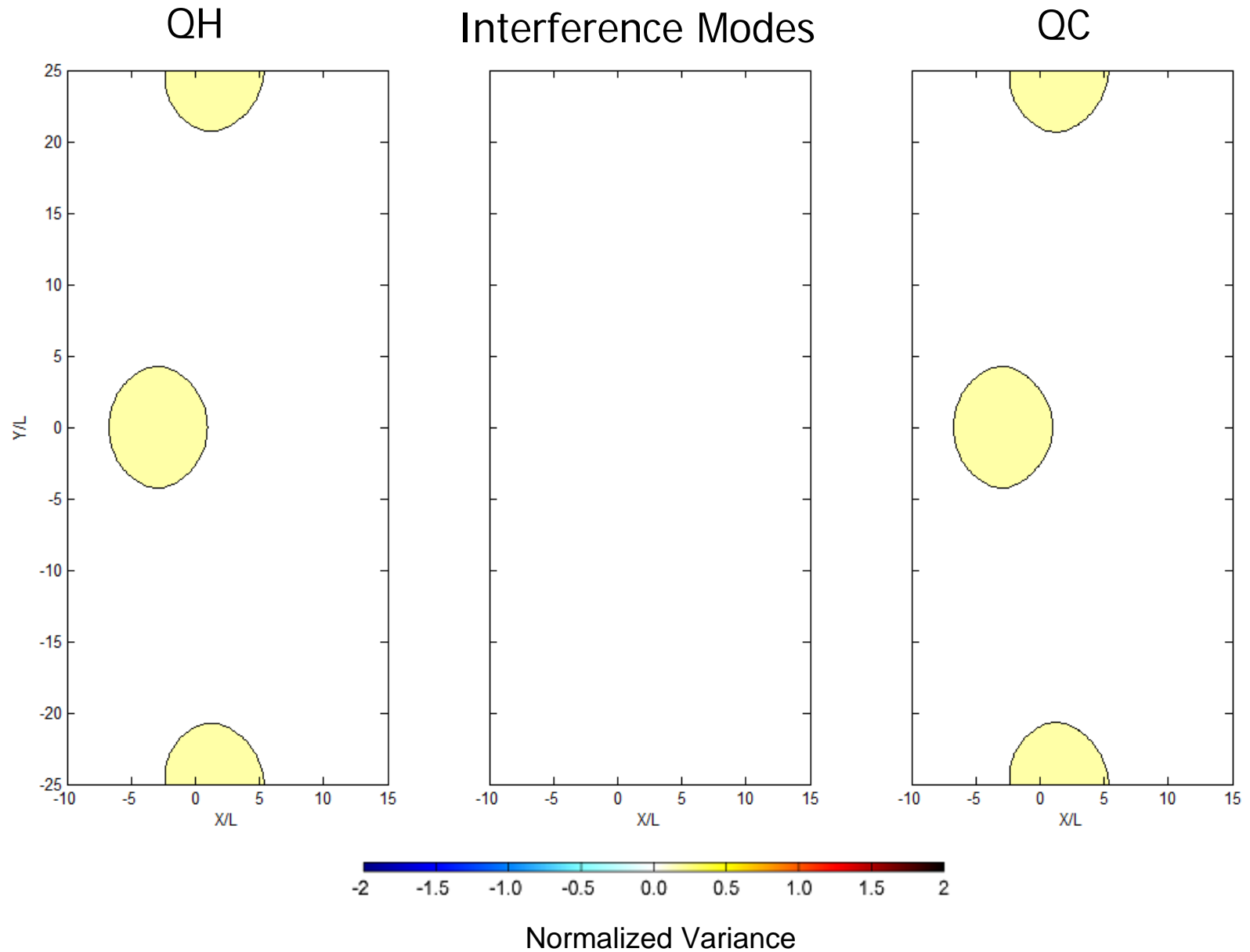
Corrections for coherent  
interference

# Example: Coherent Gaussian Packets

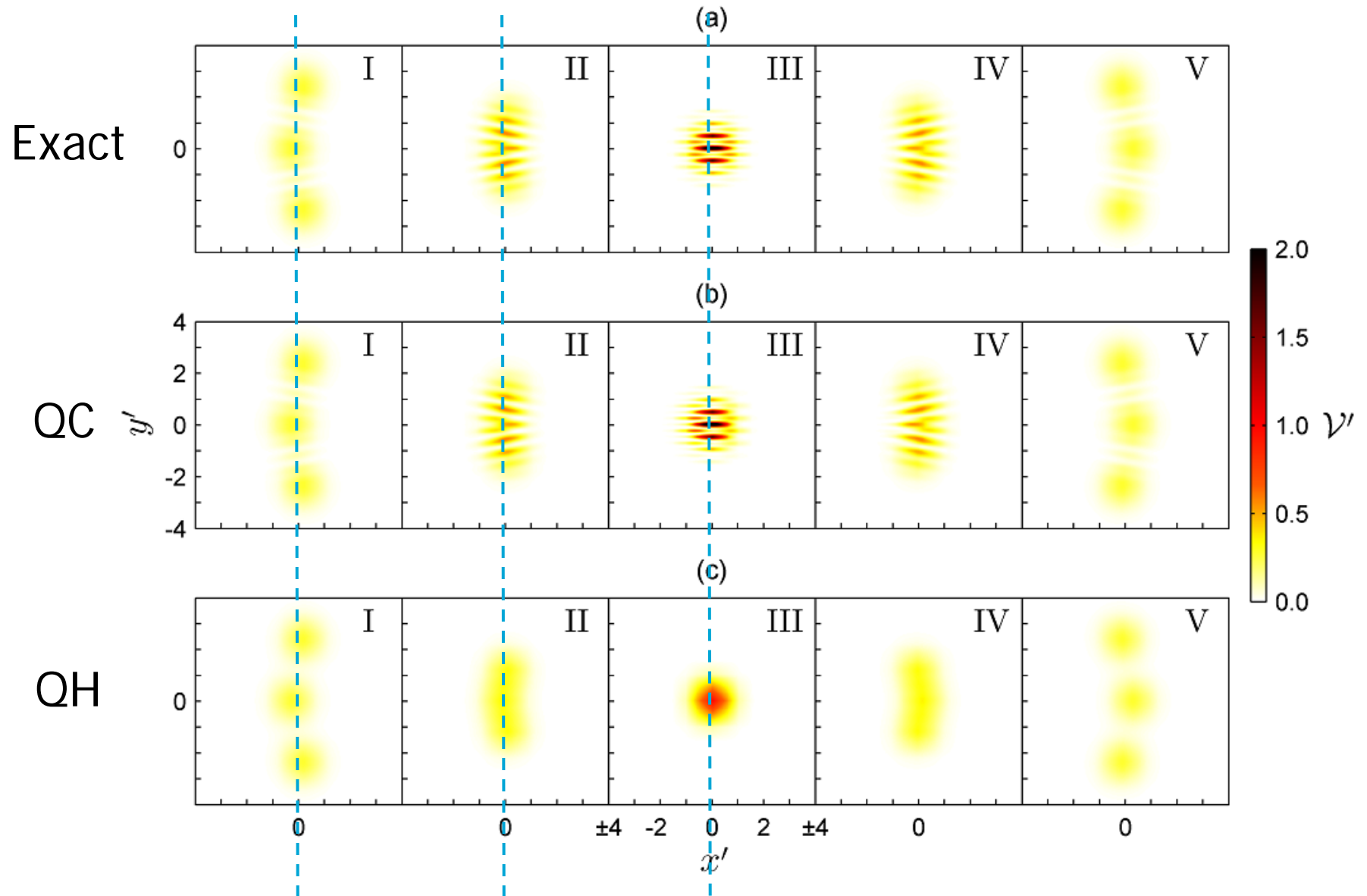


$$\zeta(\mathbf{x}, t = t') = \sum_{m=1}^N A_m \exp\left(-\alpha |\mathbf{x}|^2 + i\mathbf{k}_m \cdot \mathbf{x}\right)$$

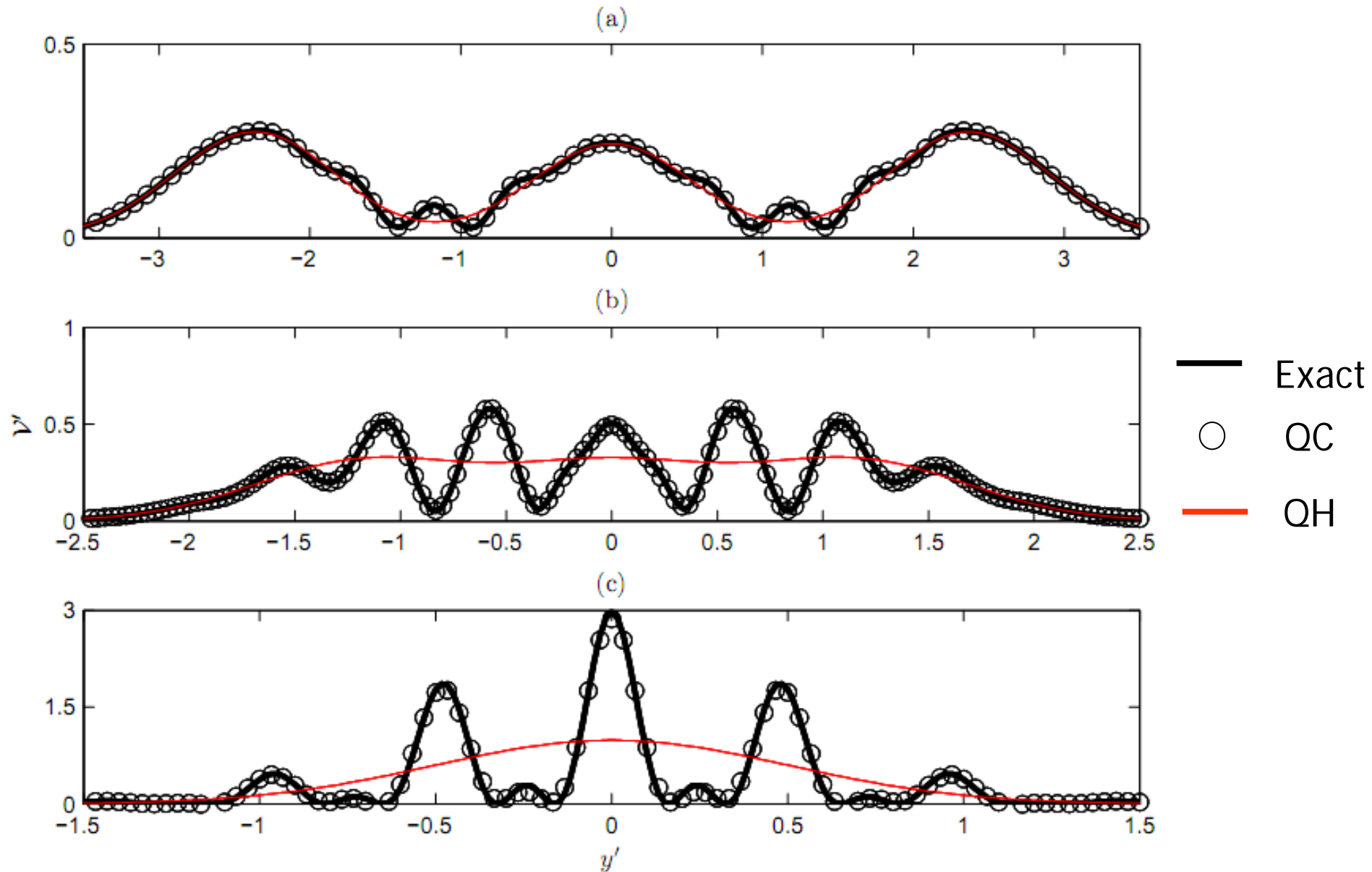
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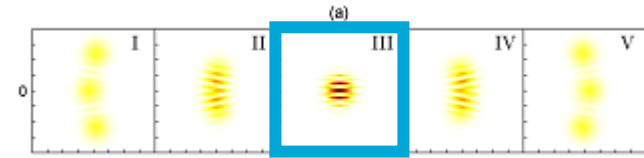
# Example: Coherent Gaussian Packets



# Example: Coherent Gaussian Packets



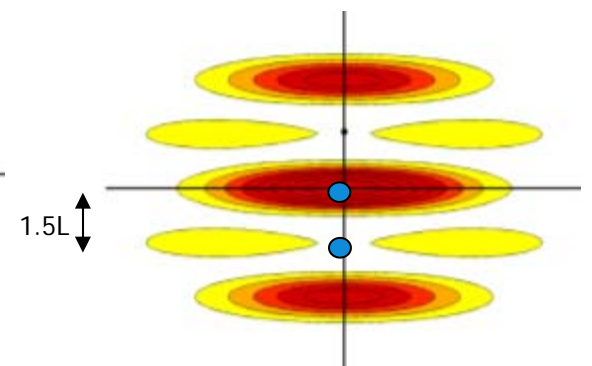
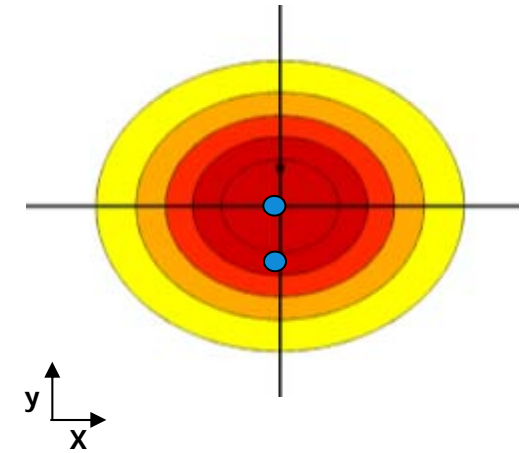
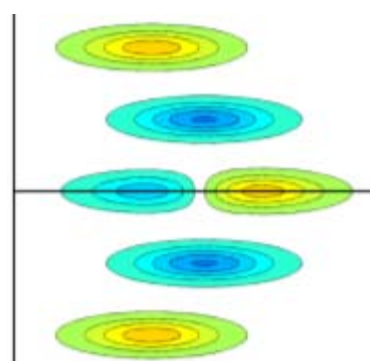
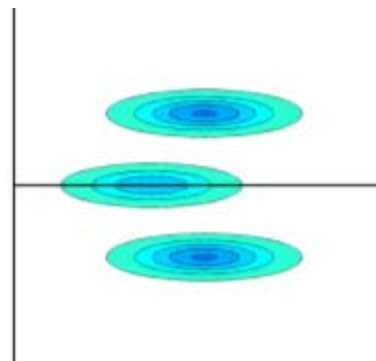
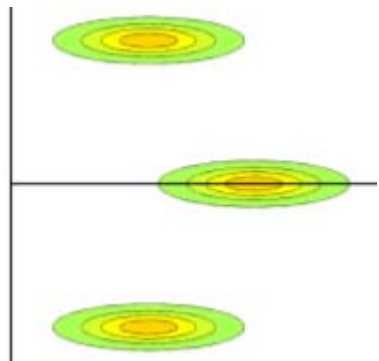
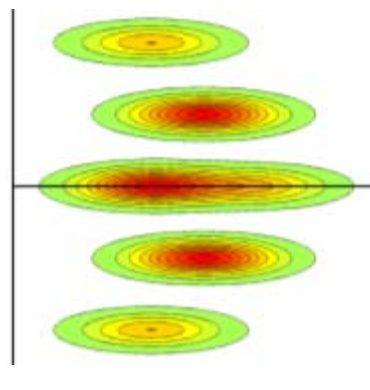
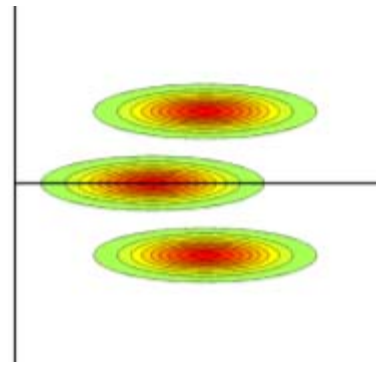
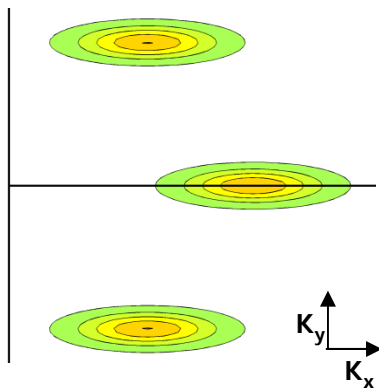
# Example: Coherent Gaussian Packets



auto-modes

Interference-modes

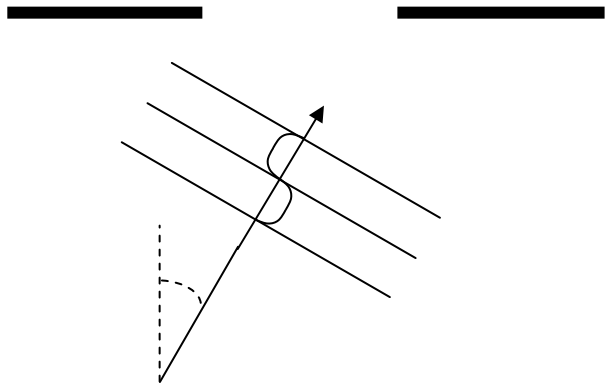
CM-spectrum





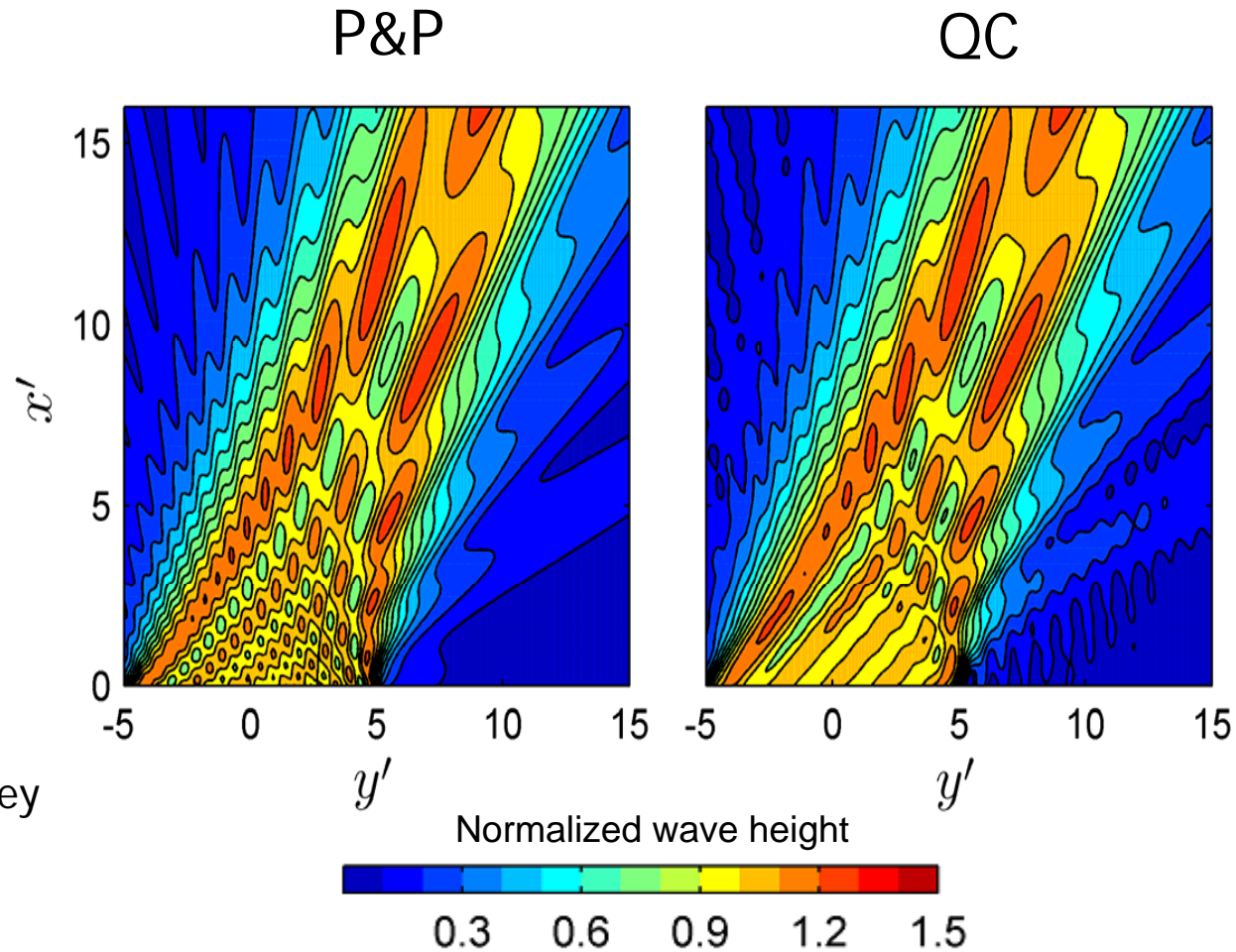
# Example: Diffraction through a gap

## Wide angle Diffraction



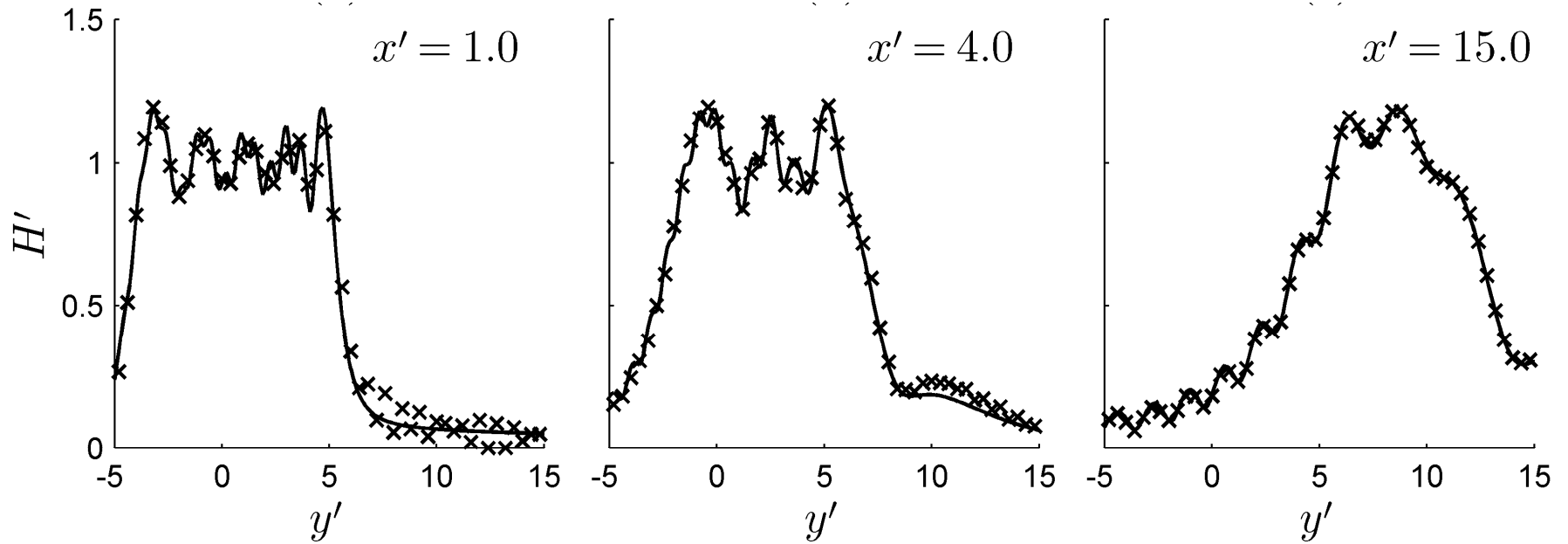
- Monochromatic wave

- Compared to solution by Penney & Price (1952)



# Examples

## Wide angle Diffraction



— P&P

× QC

# Concluding remarks

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- Consistent extension of Quasi-Homogeneous theory to inhomogeneous fields
- Compatible with operational wave modeling framework makes coupling to regional models straightforward.
- Likely areas of improvement: coastal focal zones, sheltering (headlands), and river mouths and tidal inlets.
- Application to variable topography in progress
- Same approach can be used to include non-Gaussian effects in shallow water.

